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## APPLICATION OF BLOCH NMR EQUATION AND PENNES BIOHEAT EQUATION TO THERANOSTICS

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**INTRODUCTION** Theranostics has been regarded as a key part of personalized medicine and requires considerable advances in predictive medicine; novel theranostic agents are developed and carefully designed for in vivo quantitative assessment of the amount of drug reaching a pathological region and the visualization of molecular changes due to the therapeutic effects of the delivered drug. This study intends to mathematically model a closely knitted theranostic method in which a specially selected RF field is used to heat up a tissue and at the same time cause the spins of the tissue to emit MR signals.

**MATHEMATICAL FORMULATION** We consider bioheat flow in one direction [1, 2] given in eqn(1); where  $\rho$  is tissue density,  $c$  is the specific heat of tissue,  $T$  is the tissue temperature,  $t$  is the time,  $w_b$  is the blood perfusion rate,  $c_b$  is the specific heat of blood,  $T_b$  is the supplying arterial blood temperature,  $k$  is the thermal conductivity of tissue, and  $x$  is the distance from the skin surface. SAR is the applied RF power per unit volume. If  $T$  changes very slowly with  $x$ , we have eqn(2) and solution to eqn(2) is given in eqn(3). If the  $T$  before the application of the RF field does not defer significantly from  $T_b$ , the initial the condition for this problem is given in eqn(4) and the final solution is given in eqn(5). The RF power for the voxel volume  $V_{\text{vox}}$  is  $P_{\text{rf}} = (\text{SAR}) V_{\text{vox}}$ . The energy of the oscillating radio wave is given as  $E_{\text{rf}} = (1.055 \times 10^{-34} \text{Js}) \gamma B_1$ , whose rate of change is expressed as in eqns(6) and (7). We can relate time dependent MRI signal to SAR using the time independent NMR equation [2] given by eqn (8) and (9). If we sample the signal when the  $M_y$  has the largest amplitude, we write  $M_0 \approx 0$ . Provided that the condition in eqn (10) holds, we have [2] eqn (11). From eqns (10) and (11), we obtain eqn (12). If the RF  $B_1$  field is applied at time  $t_0 = 0$ , we have eqn (13). This solution is valid for the condition in eqn (14). It is always required that the  $M_y$  be finite as time tends to infinity; therefore, the solution to the problem is given by eqn (15).

**ANALYSIS OF RESULTS** The results obtained in this study have been simulated with relaxation parameters of human liver at 1.5T [3] and the corresponding thermal properties [1, 3]:  $T_1 = 0.610\text{s}$ ,  $T_2 = 0.057\text{s}$ ,  $w_b = 2.86\text{kg/m}^3\text{s}$ ,  $c_b = 3960\text{J/kg.K}$ ,  $\rho = 1060\text{kg/m}^3$ ,  $c = 3600\text{J/kg.K}$ . Plots a and b ( $\text{SAR} = 4\text{W/m}^3$ ) give the distribution of the  $T$  and  $M_y$  on a log scale while plot c ( $\text{SAR} = 40000\text{W/m}^3$ ) gives the density plot of  $M_y$  as a function of time and tissue temperature.

**CONCLUSION** The temperature distribution and the RF power needed to generate RF  $B_1(t)$  field within the medically acceptable SAR limit during MRI scanning procedure have been investigated by solving the Pennes Bioheat equation in terms of MRI parameters. The relationship between  $T$ , SAR and RF  $B_1(t)$  at any given time is clearly shown in eqn (5), eqn (10) and Plots a, b, c.

**REFERENCES** [1] Tzu-Ching Shih, Ping Yuan, Win-Li Lin, Hong-Sen Kou. MEP 29 (2007) 946-953. [2] O. B. Awojoyogbe, M. Dada, O.P. Faromika, O.E. Dada. CMRA. Vol. 38 A (3) 85-101 (2011). [3] Bottomley PA, Foster TH, Argersinger RE, Pfeifer LM (1984). Med Phys 11:425-448.

$$\rho c \frac{\partial T}{\partial t} = k \frac{\partial^2 T}{\partial x^2} + SAR - w_b c_b (T - T_b) \quad (1) \quad T_0(t-t_0) \leq 1 \quad (14)$$

$$\rho c \frac{\partial T}{\partial t} = SAR - w_b c_b (T - T_b) \quad (2) \quad M_y(t) = C_1(\beta)^{\frac{1-T_f}{2}} (t)^{\frac{1-T_f}{2}} J_{\frac{1-T_f}{2}}\left(\frac{V_{vox}}{2h} t^2\right) \quad (15)$$

$$T(t) = T_b + \frac{SAR}{w_b c_b} + A \exp\left(-\frac{w_b c_b}{\rho c} t\right) \quad (3)$$

$$T(t=0) = T_b \quad (4)$$

$$T(t) = T_b + \frac{SAR}{w_b c_b} \left\{ 1 - \exp\left(-\frac{w_b c_b}{\rho c} t\right) \right\} \quad (5)$$

$$P_{rf} = \frac{dE_{rf}}{dt} \text{ and } E_{rf} = \int_{t_0}^t (SAR) V_{vox} dt \quad (6)$$

$$\gamma B_1(t) = \frac{V_{vox}}{h} (SAR)(t - t_0) \quad (7)$$

$$\frac{d^2 M_y}{dt^2} + T_0 \frac{dM_y}{dt} + (T_x + \gamma^2 B_1^2(t)) M_y = \frac{M_0}{T_1} \gamma B_1(t) \quad (8)$$

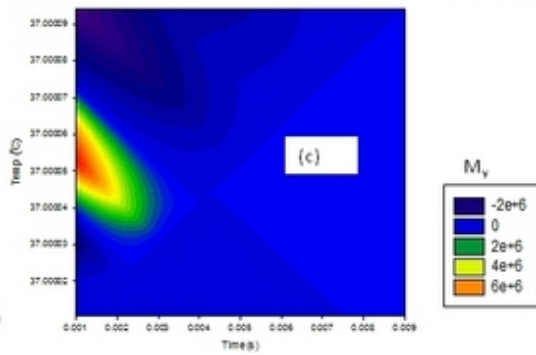
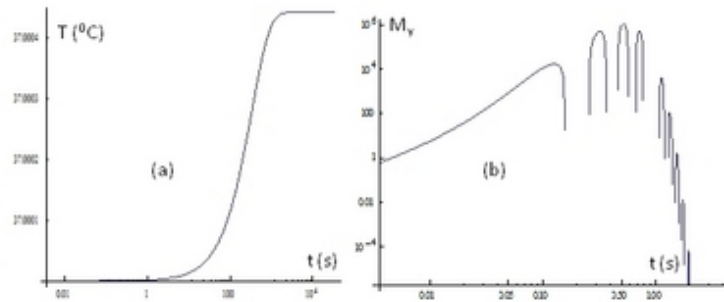
$$\text{where } T_x = \frac{1}{T_1 T_2} \text{ and } T_0 = \frac{1}{T_1} + \frac{1}{T_2} \quad (9)$$

$$T_x < \gamma^2 B_1^2(t) \quad (10)$$

$$\frac{d^2 M_y}{dt^2} + T_0 \frac{dM_y}{dt} + \gamma^2 B_1^2(t) M_y = 0 \quad (11)$$

$$\frac{d^2 M_y}{dt^2} + T_0 \frac{dM_y}{dt} + \left(\frac{V_{vox}}{h}\right)^2 (SAR)^2 (t-t_0)^2 M_y = 0 \quad (12)$$

$$M_y(t) = (\beta)^n(t) \left[ C_1 J_n\left(\frac{V_{vox}}{2h} t^2\right) - C_2 Y_n\left(\frac{V_{vox}}{2h} t^2\right) \right] \quad n = \frac{1-T_f}{2} \quad (13)$$



Disclosure of author financial interest or relationships:

**M.O. Dada**, None; **B.O. Awojogbe**, None; **S. Baroni**, None; **M.A. Aweda**, None.